

PHYSICAL MECHANISMS OF INTERFACE-MEDIATED INTERVALLEY COUPLING IN SILICON

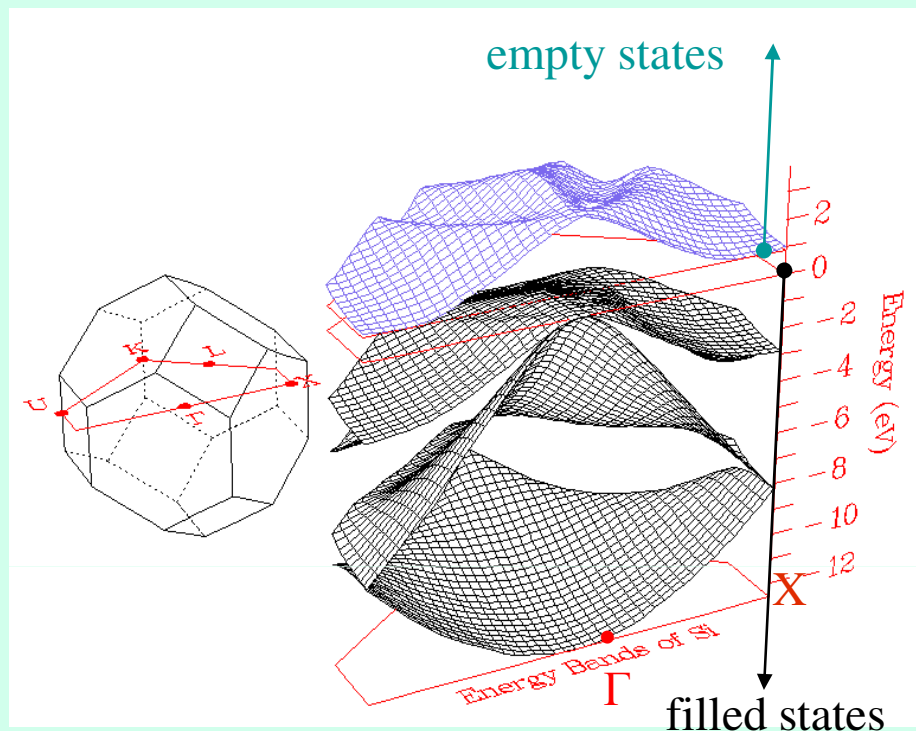
André L. Saraiva – UFRJ – Brazil



Si Qubit Workshop – Berkeley 2009

Saraiva *et al.* PRB **80**, 081305R (2009)

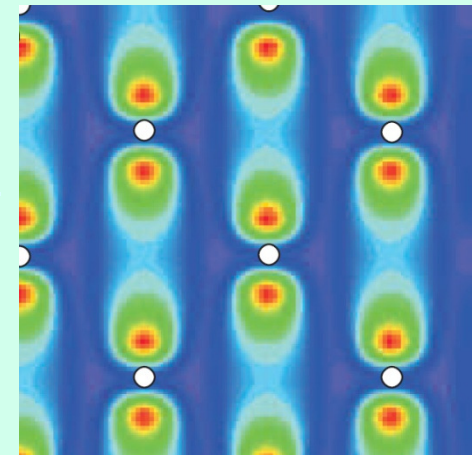
Bulk Si Electronic Structure



Bloch States

$$\phi_{k_\mu}(\vec{r}) = e^{i\vec{k}_\mu \cdot \vec{r}} u_{\vec{k}_\mu}(\vec{r})$$

$$|\phi_{k_\mu}(\vec{r})|^2 \rightarrow$$

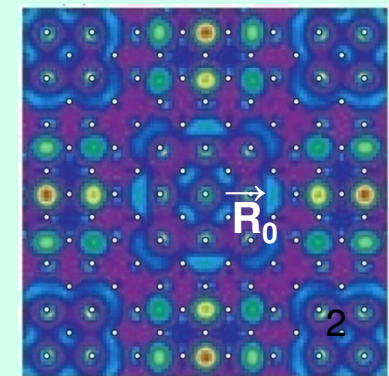


Koiller *et al* [PRB(2004)]

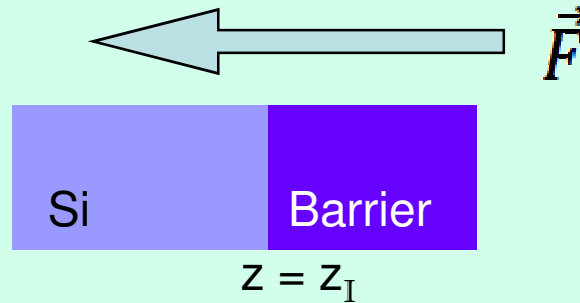
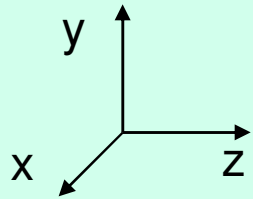
Any superposition of degenerate Bloch states is also an eigenstate (not in Bloch's form).

Example:

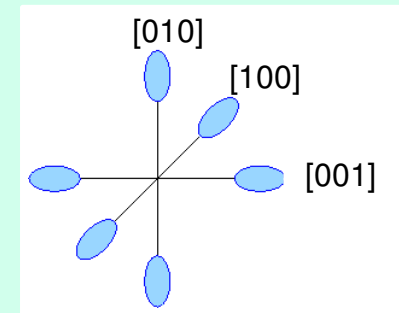
$$\phi(\vec{r}) = \frac{1}{\sqrt{6}} \sum_{\mu=1,6} \phi_{\vec{k}_\mu}(\vec{r} - \vec{R}_0) \Rightarrow |\phi(\vec{r})|^2 \quad \mapsto$$



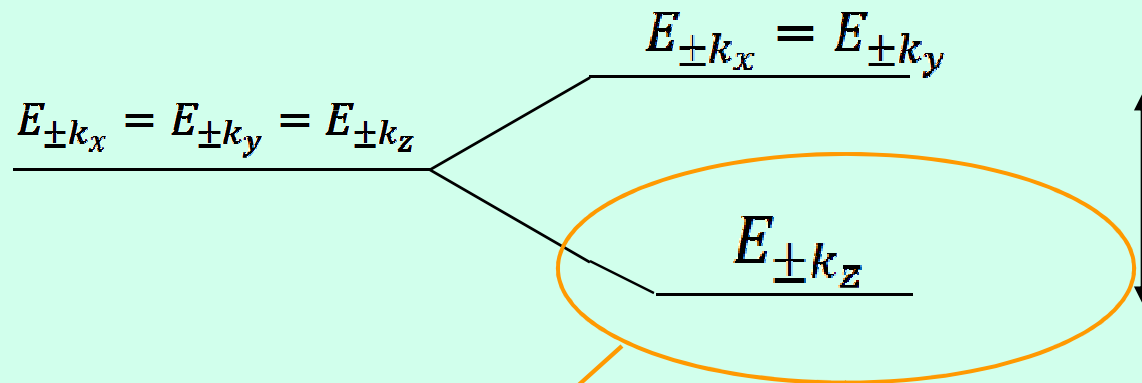
Lifting the Degeneracy



$$k_0 \approx 0.85 \frac{2\pi}{a_0}$$



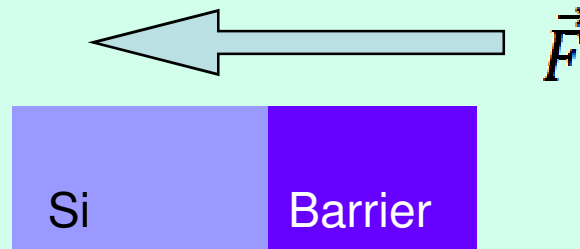
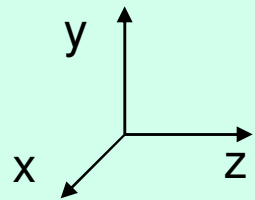
$$H = H_0 + H'$$



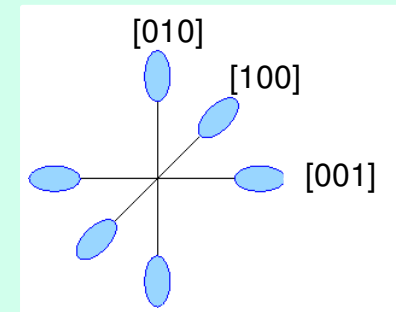
$$E_0 \propto \left(\frac{F^2}{m_z} \right)^{\frac{1}{3}} \xrightarrow{\text{yields}} \Delta_{xz} \approx 20 - 30 \text{ meV}$$

$$\varphi(\vec{r}) = a_+ \phi_{+z} + a_- \phi_{-z}$$

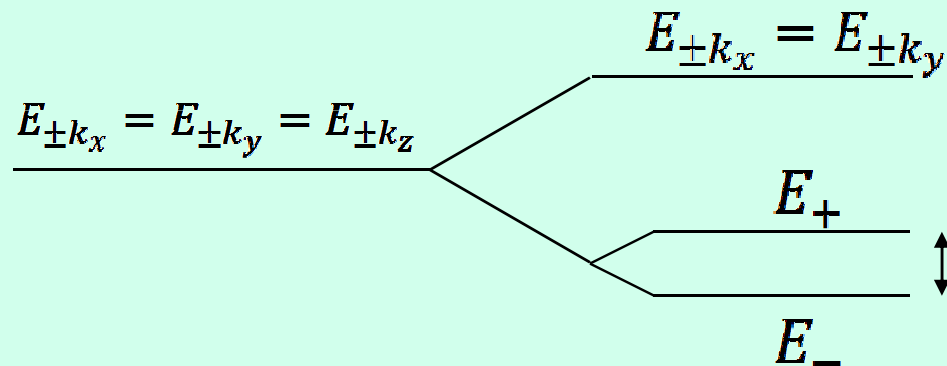
Lifting the Degeneracy



$$k_0 \approx 0.85 \frac{2\pi}{a_0}$$



$$H = H_0 + H'$$



$$2|V_{VO}|$$

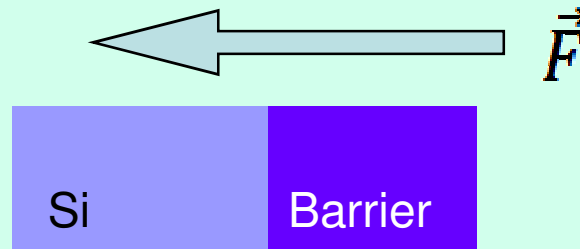
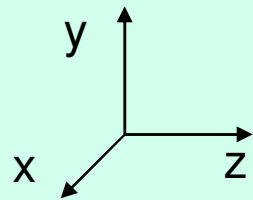
$$\approx 0.01 - 1 \text{ meV}$$

$$\approx 20 - 30 \text{ meV}$$

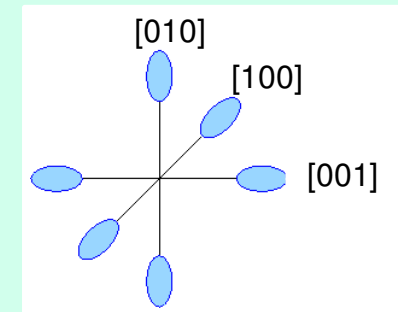
$$H = \begin{pmatrix} E_0 & V_{VO} \\ V_{VO}^* & E_0 \end{pmatrix}$$

$$V_{VO} = \langle \phi_- | H' | \phi_+ \rangle$$

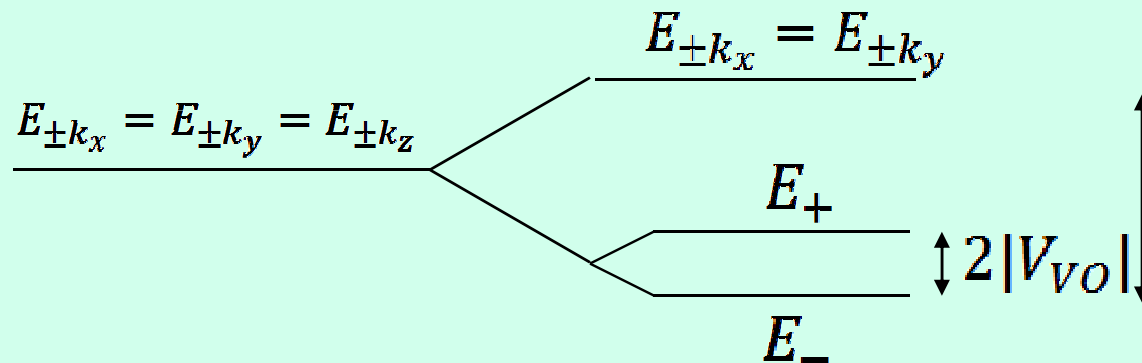
Lifting the Degeneracy



$$k_0 \approx 0.85 \frac{2\pi}{a_0}$$



$$H = H_0 + H'$$



$$E_0 \propto \left(\frac{F^2}{m_z} \right)^{\frac{1}{3}} \xrightarrow{\text{yields}} \Delta_{xz} \approx 20 - 30 \text{ meV}$$

PRL 96, 236801 (2006)

PHYSICAL REVIEW LETTERS

week ending
16 JUNE 2006

Valley Polarization in Si(100) at Zero Magnetic Field

$2|V_{VO}| \rightarrow 23 \text{ meV}$

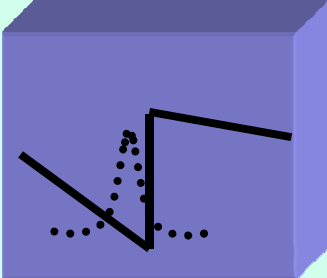
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Effective Mass Theory – Step Model

$$H' = U(z) - eFz \Rightarrow$$


$$U(z) = U_0 \Theta(z - z_I)$$

$z = z_I$

- Single Valley Eff. Mass Equation

Finite Differences Method
(Fine Description of Evanescent Tail)

$$\left[-\frac{\hbar^2}{2m_z} \frac{\partial^2}{\partial z^2} + U(z) - eFz \right] \psi(z) = E_0 \psi(z)$$

$$\phi_{\pm}(\vec{r}) = \psi(z) e^{\pm i \vec{k}_0 \cdot \vec{r}} \left[\sum_{\vec{G}} c_{\pm}(\vec{G}) e^{i \vec{G} \cdot \vec{r}} \right]$$

$$V_{VO} = \langle \phi_- | H' | \phi_+ \rangle$$

Plane Waves Expansion Coefficients
Kohn Orbitals- DFT

Effective Mass Theory – Step Model

$$z_I = 0 \quad V_{VO} = \sum_{\vec{G}, \vec{G}'} c_+^*(\vec{G}) c_-^*(\vec{G}') \delta(G_x - G_x') \delta(G_y - G_y') \times I(G_z, G_z')$$

$$I(G_z, G_z') = \int_{-\infty}^{\infty} |\Psi(z)|^2 e^{iQz} [U_0 \Theta(z)] dz - \int_{-\infty}^{\infty} |\Psi(z)|^2 e^{iQz} [eFz] dz$$

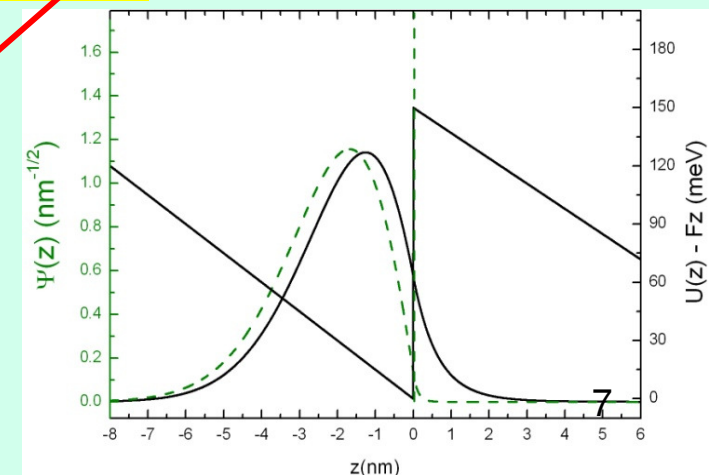
$$= \int_{-\infty}^{\infty} \frac{i e^{iQz}}{Q} |\Psi(z)|^2 U_0 \delta(z) dz + \int_0^{\infty} \frac{i e^{iQz}}{Q} \frac{d}{dz} |\Psi(z)|^2 U_0 dz + I_F$$

$$= i \frac{U_0}{Q} |\Psi(0)|^2 + i \frac{U_0}{Q} \int_0^{\infty} \frac{d}{dz} |\Psi(z)|^2 e^{iQz} dz + I_F$$

$$Q \stackrel{\text{def}}{=} G_z - G_z' - 2k_0$$

$$V_{VO} = V_{\delta} + V_T + V_F$$

~ 0

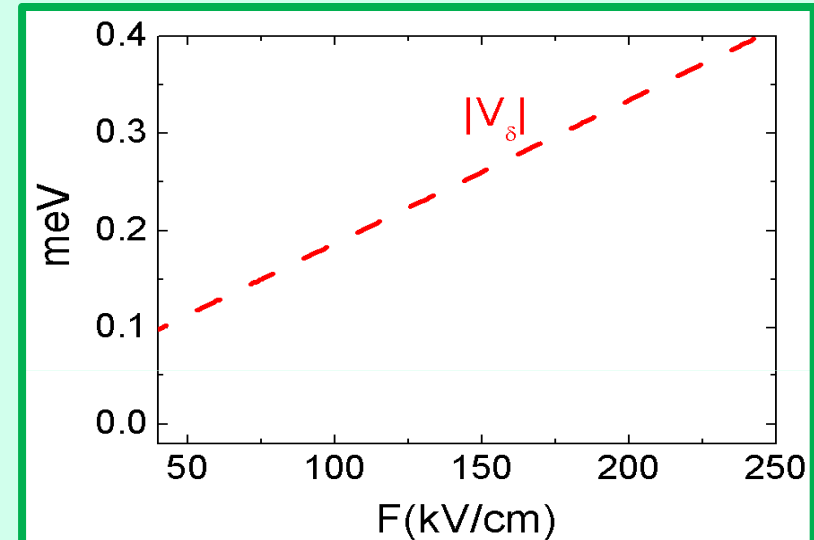
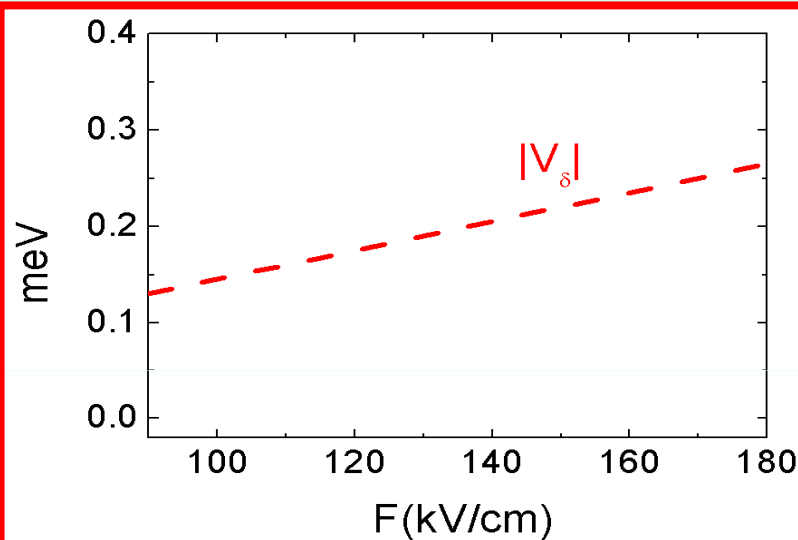


Step Model Results

$$V_{\delta} = U_0 |\Psi(0)|^2 \sum c_- c_+^* \delta_x \delta_y \frac{i}{Q}$$

$U_0 = 150 \text{ meV} \sim \text{SiGe}$

$U_0 = 3 \text{ eV} \sim \text{SiO}_2$

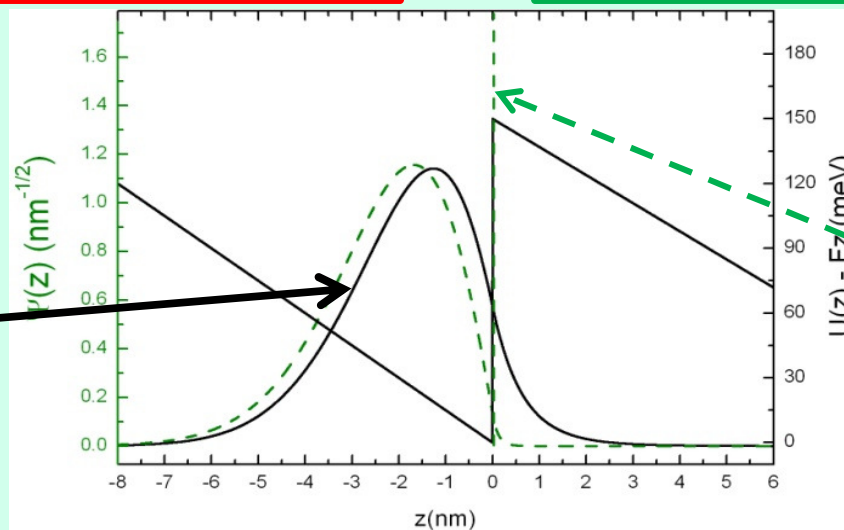


Small U_0
Vs
Large $|\Psi(0)|^2$

Large U_0
Vs
Small $|\Psi(0)|^2$

SiGe

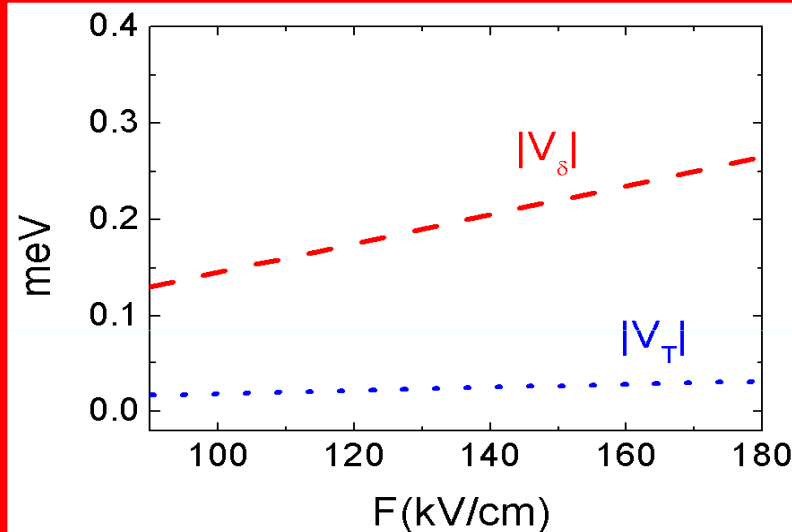
SiO₂



Step Model Results

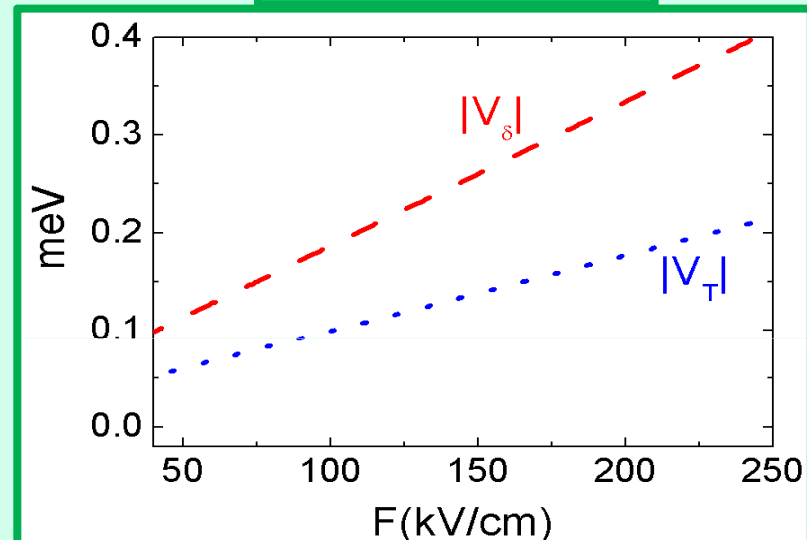
$$V_T = \sum c_- c_+^* \delta_x \delta_y \int_0^\infty \frac{i e^{iQz}}{Q} \frac{d}{dz} |\Psi(z)|^2 U_0 dz$$

$U_0 = 150 \text{ meV} \sim \text{SiGe}$



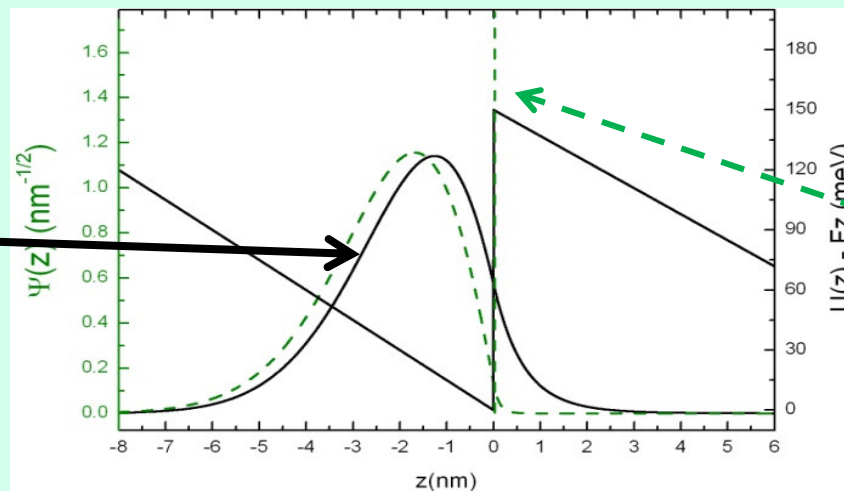
Evanescent Tail much larger than Wavelength

$U_0 = 3 \text{ eV} \sim \text{SiO}_2$



Evanescent Tail comparable to Wavelength

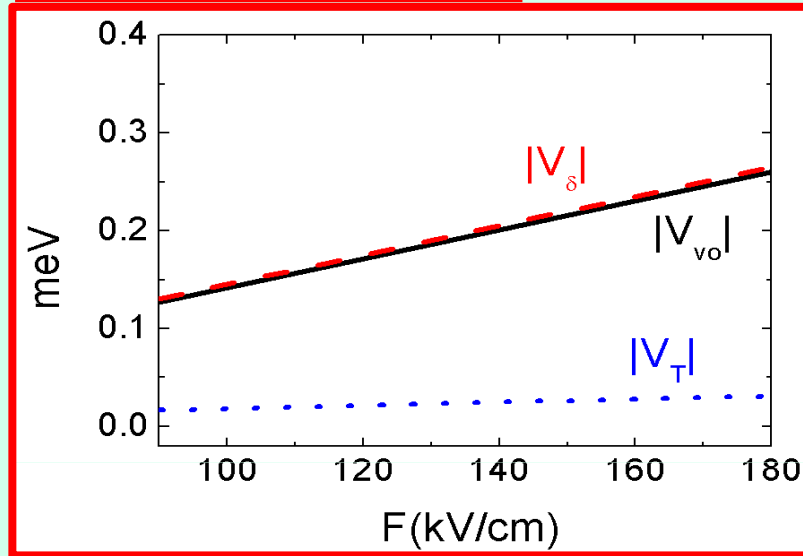
SiGe



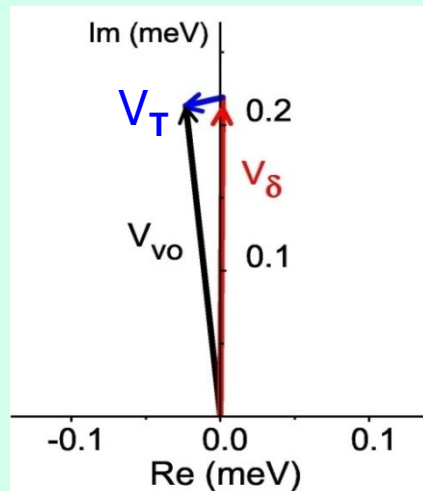
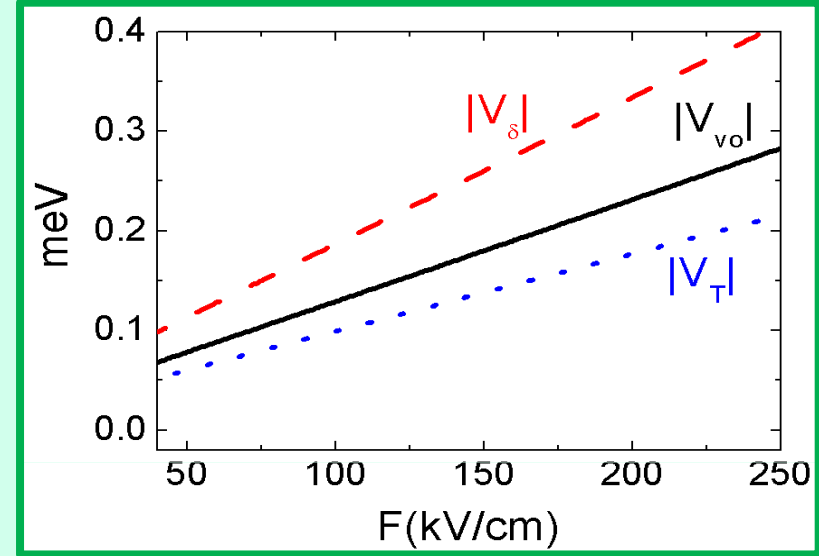
SiO₂

Step Model Results

$U_0 = 150 \text{ meV} \sim \text{SiGe}$



$U_0 = 3 \text{ eV} \sim \text{SiO}_2$

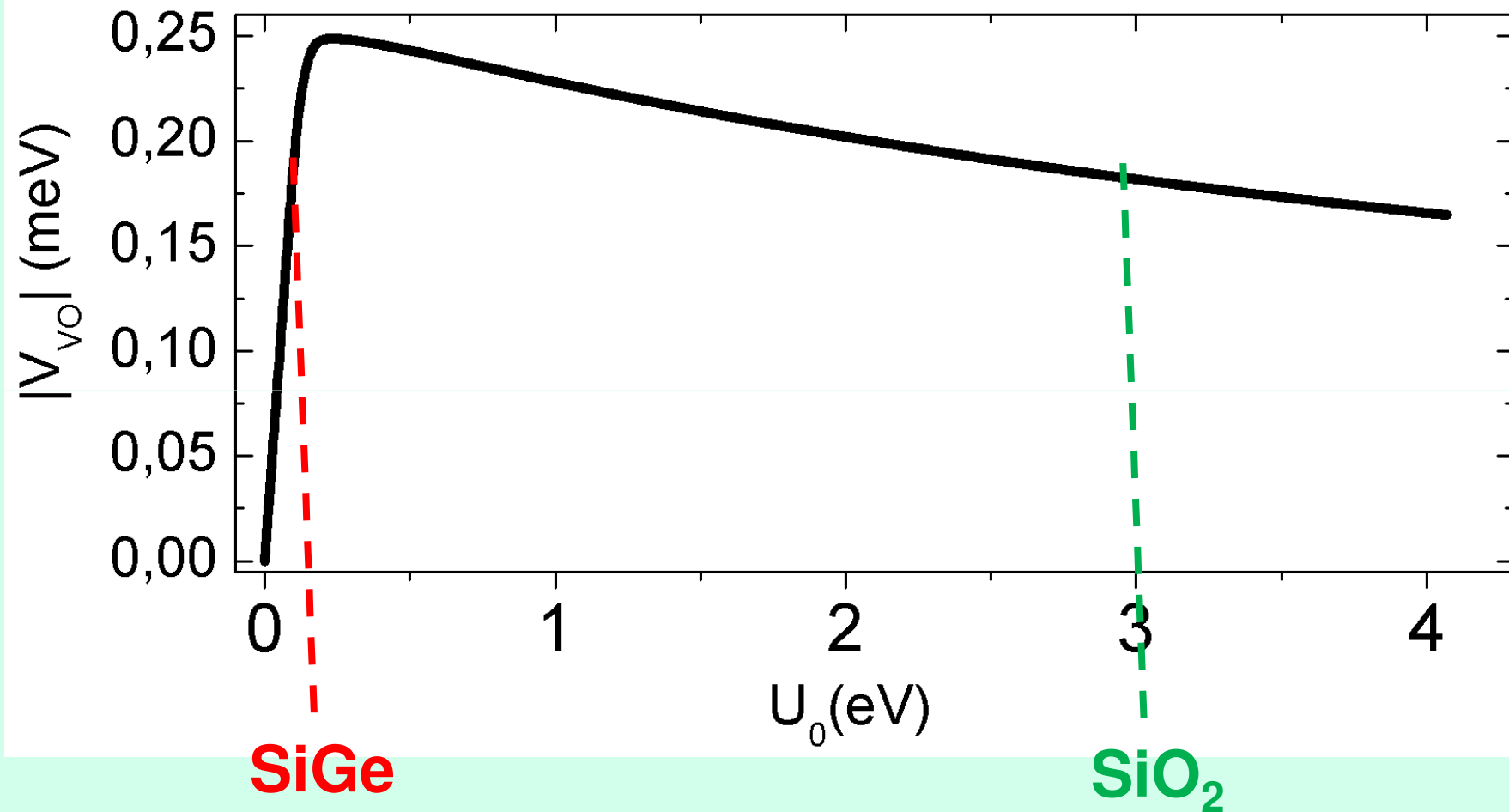


$$V_{VO} = V_\delta + V_T + V_F$$

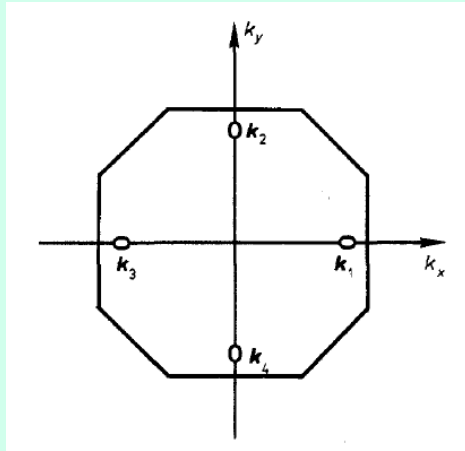
$$|V_{VO}| \approx 10 \text{ meV}$$

$$F \approx 10\,000 \text{ kV/cm!}$$

Role of the Barrier Material



Relevance of Umklapp Processes



R. Resta,
J. Phys. C (1977)

$$V_\delta = v |\Psi(0)|^2$$

$$\phi_\pm(\vec{r}) = \psi(z) e^{\pm i \vec{k}_0 \cdot \vec{r}} \left[\sum_{\vec{G}} c_\pm(\vec{G}) e^{i \vec{G} \cdot \vec{r}} \right]$$

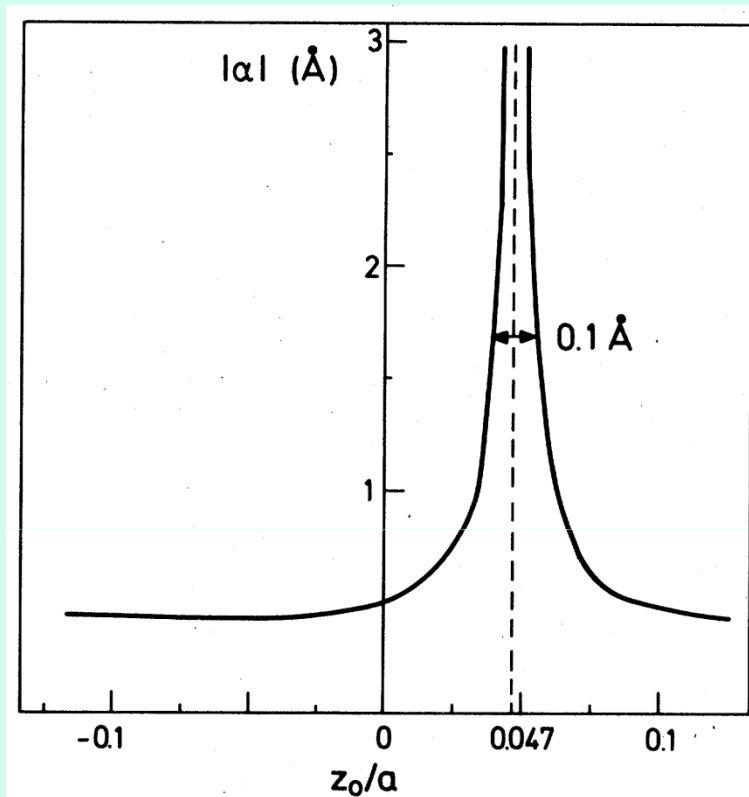
Taking only $\mathbf{G}=\mathbf{G}'=0$

$$v \approx 5.1 \times 10^{-11} U_0 \text{ (eV} \cdot \text{m)}$$

**Convergence on G
(up to second nearest neighbors from Γ in the BCC RL)**

$$v \approx 1.6 \times 10^{-11} U_0 \text{ (eV} \cdot \text{m)}$$

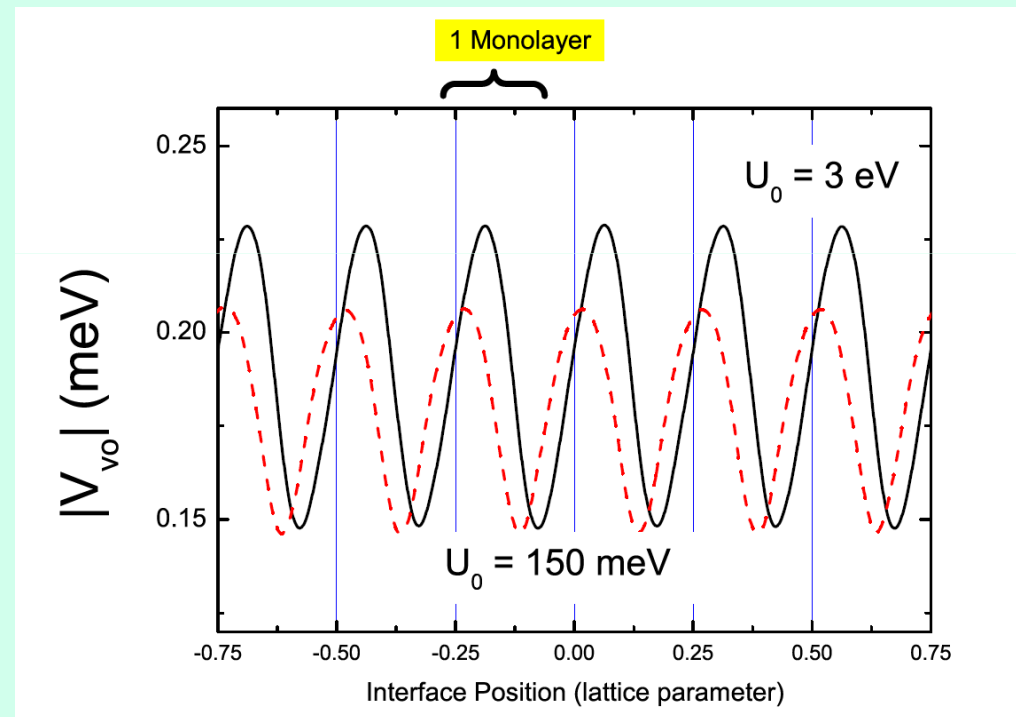
Interface Position



← 1 monolayer →

Scattering Length
L. Sham & M Nakayama,
PRB(1979)

$$H = H_0 + U_0 \Theta(z - z_I) - eFz$$

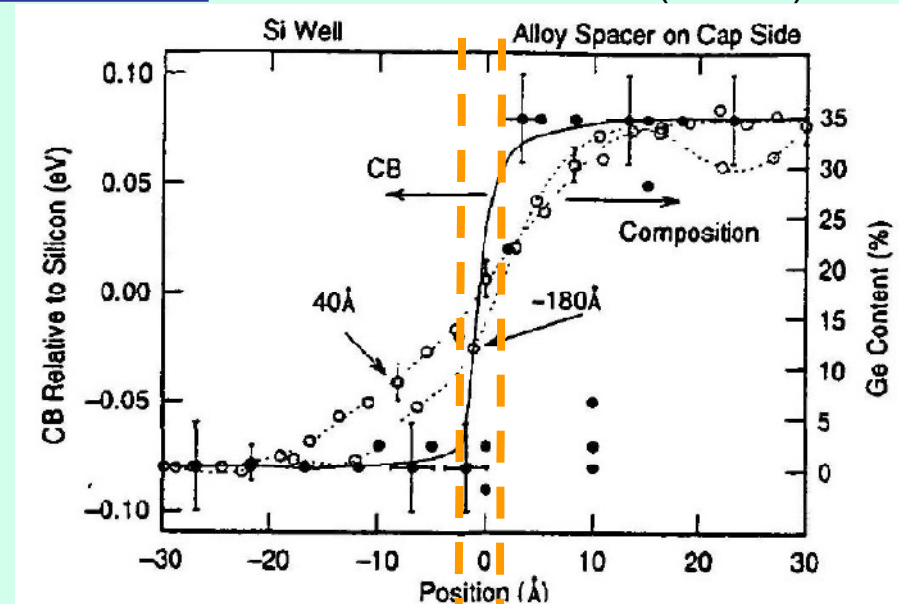
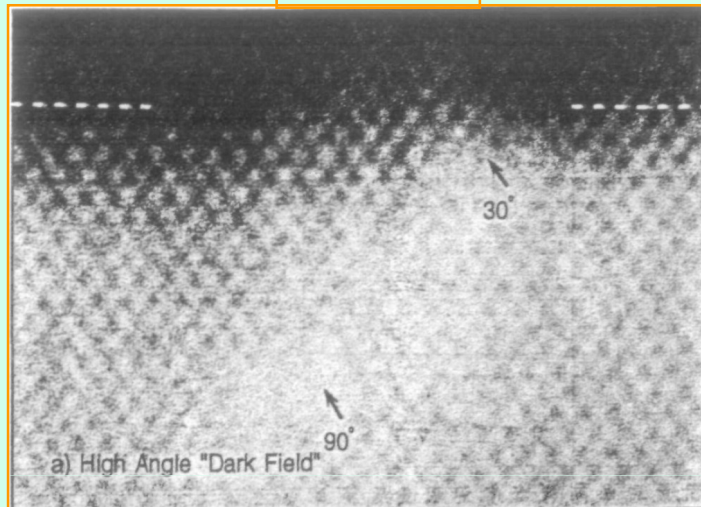


(Present Work)

Real Samples

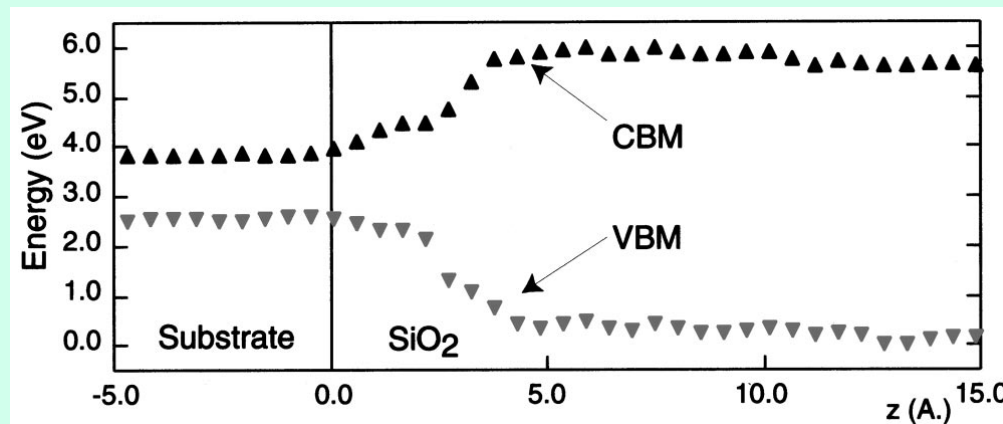
(EELS)

Si/SiGe



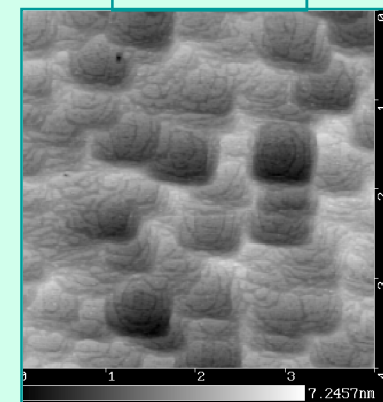
Batson [J. Elec. Micr.(2000)]

≈ 1 – 2 ML



Yamasaki *et al* [PRB(2004)]

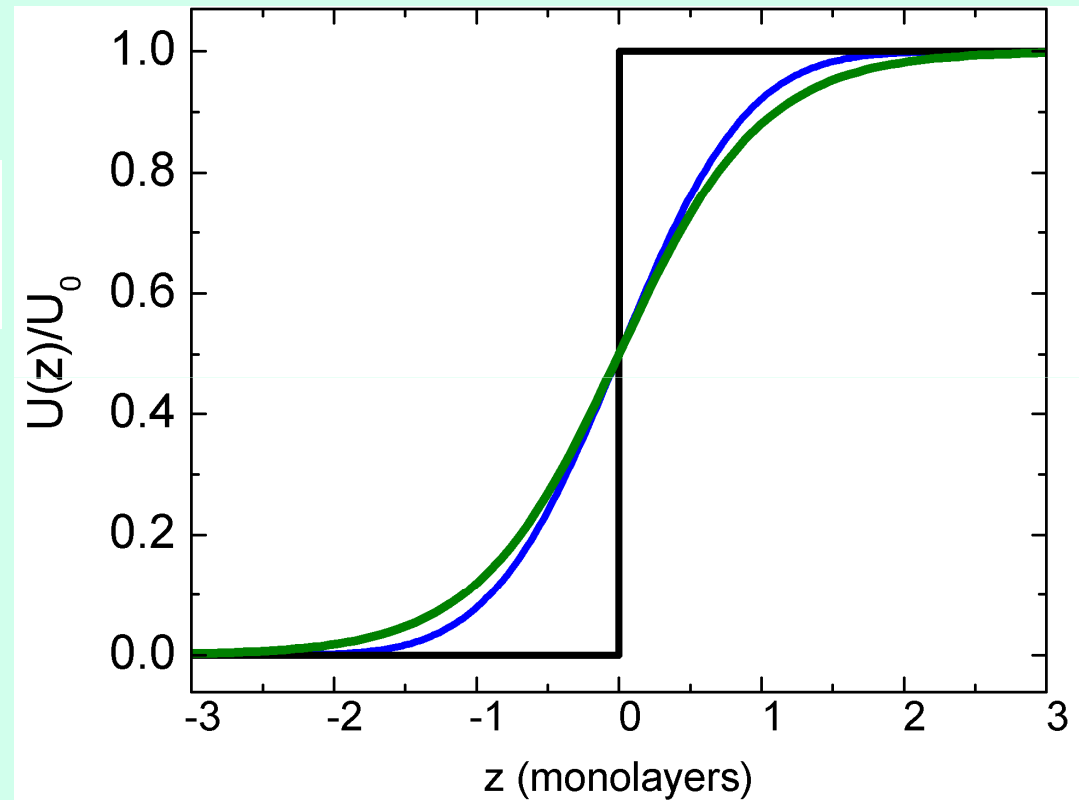
Si/SiO₂



Finite Width Models

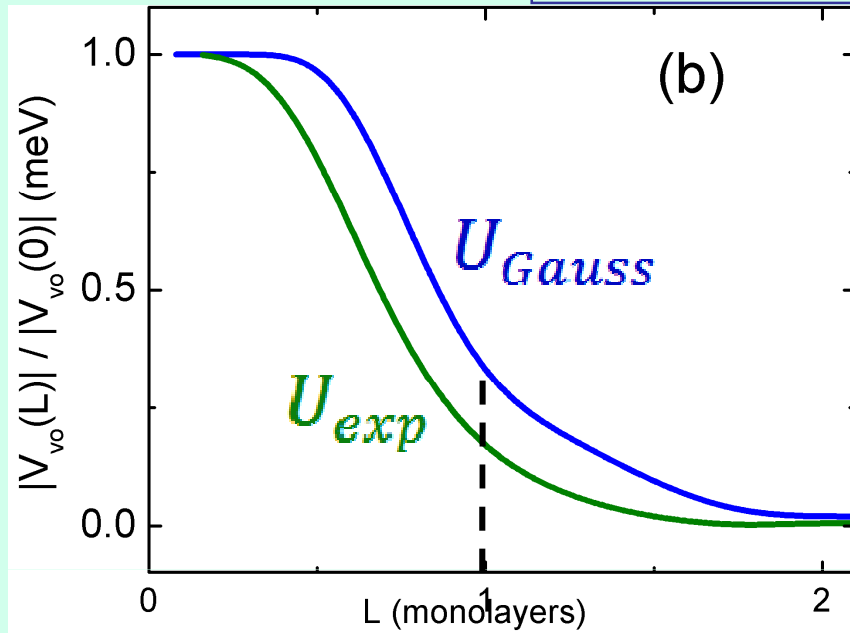
$$U_{exp} = \frac{U_0}{2} \left[\tanh\left(\frac{z}{L}\right) + 1 \right]$$

$$U_{Gauss} = \frac{U_0}{2} \operatorname{erfc}\left(-\frac{z}{L}\right)$$



RMS width of U_{exp} is 8% larger than
RMS width of U_{Gauss}

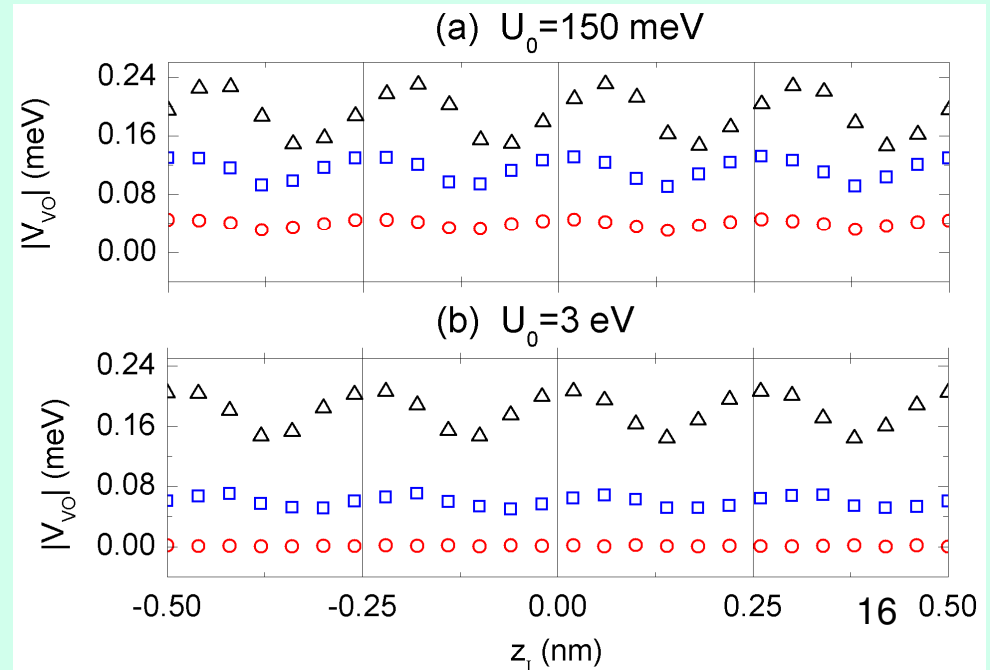
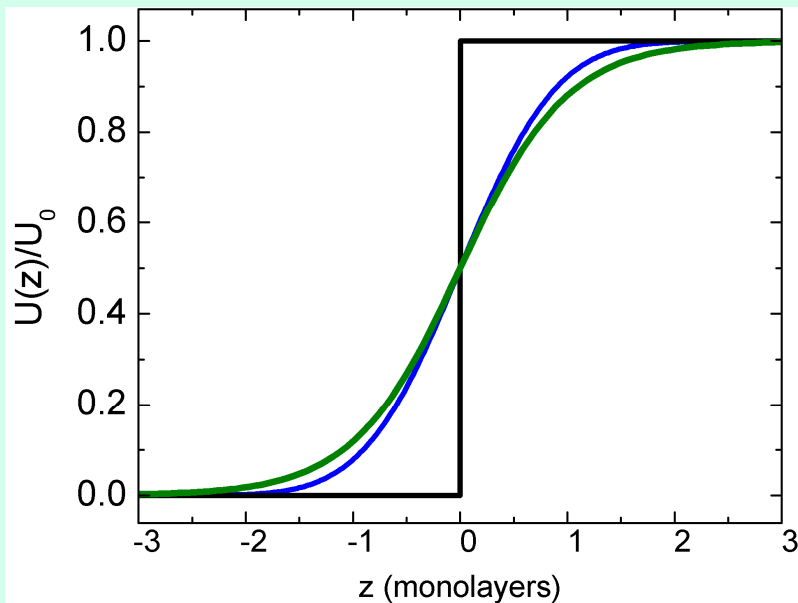
Results for Finite Width



RMS width of U_{exp} is 8% larger than RMS width of U_{Gauss}

V_{VO} of U_{Gauss} is almost twice the V_{VO} of U_{exp}

L = 1 ML



Conclusions

- We identified nanofabrication parameters that might enhance interface-induced intervalley coupling;
- High quality heterostructures ($L \rightarrow 0$) are indeed more suitable for QC; no sign of intervalley coupling as large as 10 meV;
- V_{VO} could be further enhanced by some band structure engineering (design of suitable U_0);
- Effective Mass theory gives consistent description for the intervalley coupling, free of fitting parameters;
- The unphysical oscillations of V_{VO} with the interface position are artifacts of the effective mass theory in combination with atomistic Bloch Functions;



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Sankar Das Sarma – Maryland

Belita Koiller - UFRJ

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